

Title: Re-examination of Stieltjes Constants

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July 3, 2026

The Stieltjes Constants are typically defined as:

$$\gamma_n = \lim_{m \rightarrow \infty} \sum_{k=1}^m \frac{\ln^n(k)}{k} - \frac{\ln^{n+1}(m)}{n+1} \quad (1)$$

In the last Stieltjes post, I found the following relation:

$$\gamma_n = \lim_{m \rightarrow \infty} \sum_{q=0}^n \frac{\binom{n}{q}}{n!} (-1)^{n-q} q^n \int_0^1 dt \frac{1-t^m}{1-t} \lim_{z \rightarrow 0} \partial_z^n \left(\frac{(-\ln(t))^{-z}}{\Gamma(1-z)} \right) - \frac{\ln^{n+1}(m)}{n+1} \quad (2)$$

Evaluating the sum in q :

$$\sum_{q=0}^n \frac{\binom{n}{q}}{n!} (-1)^{n-q} q^n = 1$$

Making the second term an integral representation:

$$\frac{\ln^{n+1}(m)}{n+1} = (-1)^{n+1} \frac{\ln^{n+1}\left(\frac{1}{m}\right)}{n+1} = (-1)^n \int_{\frac{1}{m}}^1 \frac{\ln^n(t)}{t}$$

To calculate the n th derivative in z , we calculate:

$$\begin{aligned} \left(\partial_z^p ((-\ln(t))^{-z}) \right)_{z \rightarrow 0} &= (-\ln(-\ln(t)))^p \\ \left(\partial_z^r \left(\frac{1}{\Gamma(1-z)} \right) \right)_{z \rightarrow 0} &= B_r \left(\psi_0(1), -\psi_1(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right) \end{aligned} \quad (3)$$

where $B_r(x_0, x_1, \dots, x_{n-1})$ is the complete exponential Bell polynomials:

$$B_r \left(\psi_0(1), -\psi_1(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right) = \sum_{\nu=1}^r B_{r,\nu} \left(\psi_0(1), -\psi_1(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right)$$

where $B_{n,k}(x_0, x_1, \dots, x_{n-1})$ are the incomplete Bell polynomials. The complete exponential Bell polynomials have the following generating function:

$$\exp \left(\sum_{a=1}^{\infty} x_a \frac{t^a}{a!} \right) = 1 + \sum_{b=1}^{\infty} B_b(x_1, x_2, \dots, x_b) \frac{t^b}{b!}$$

So, the n th derivative in Equation 2 is:

$$\begin{aligned}
\lim_{z \rightarrow 0} \partial_z^n \left(\frac{(-\ln(t))^{-z}}{\Gamma(1-z)} \right) &= \lim_{z \rightarrow 0} \sum_{\omega=0}^n \binom{n}{\omega} \partial_z^{n-\omega} ((-\ln(t))^{-z}) \partial_z^\omega \left(\frac{1}{\Gamma(1-z)} \right) \\
&= \sum_{\omega=0}^n \binom{n}{\omega} (-\ln(-\ln(t)))^{n-\omega} B_\omega \left((-1)^0 \psi_0(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right)
\end{aligned} \tag{4}$$

All of this work leads to the following restatement of Equation 2:

$$\begin{aligned}
\gamma_n &= \lim_{m \rightarrow \infty} - \int_{\frac{1}{m}}^1 dt \frac{(-\ln(t))^n}{t} + \int_0^1 dt \\
&\quad \sum_{\omega=0}^n \left(\binom{n}{\omega} \frac{1-t^m}{1-t} (-\ln(-\ln(t)))^{n-\omega} B_\omega \left((-1)^0 \psi_0(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right) \right)
\end{aligned} \tag{5}$$

Now, because we have two terms going to ∞ that we subtract to give a finite answer, we cannot simply use L'Hopital's Rule with a ratio of those terms because the finite answer part will be lost during differentiation. The most advisable path is to series expand the two terms in m , cancel the divergent parts, and see what remains. Taking a different approach to my previous post, we find that:

$$\begin{aligned}
\frac{\ln^{n+1}(m)}{n+1} &= \int_{\frac{1}{m}}^1 dt \frac{(-\ln(t))^n}{t} \\
&= (-1)^{n-1} n \int_{\frac{1}{m}}^1 dt \int_t^1 dt' \frac{(\ln(t'))^{n-1}}{t't} \\
&= n! \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-1} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) (-\ln(t_{n-1})) \\
&= n! \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-1} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) \sum_{s=1}^{\infty} \frac{(1-t_{n-1})^s}{s} \\
&= n! \sum_{s=1}^{\infty} \frac{1}{s} \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-2} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) \int_0^{1-t_{n-2}} dt_{n-1} \frac{t_{n-1}^s}{1-t_{n-1}} \\
&= n! \sum_{s=1}^{\infty} \frac{1}{s} \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-2} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) B_{1-t_{n-2}}(s+1, 0) \\
&= n! \sum_{s=1}^{\infty} \frac{1}{s} \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-3} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) \sum_{w_1=0}^{\infty} \frac{1}{w_1+s+1} \int_0^{1-t_{n-3}} dt_{n-2} \frac{t_{n-2}^{w_1+s+1}}{1-t_{n-2}} \\
&= n! \sum_{s=1}^{\infty} \frac{1}{s} \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-3} \left(\int_{t_{\mu-1}}^1 dt_\mu \frac{1}{t_\mu} \right) \sum_{w_1=0}^{\infty} \frac{1}{w_1+s+1} \sum_{w_2=0}^{\infty} \frac{(1-t_{n-3})^{w_2+w_1+s+2}}{w_2+w_1+s+2} \\
&= n! \sum_{s=1}^{\infty} \frac{1}{s} \int_{\frac{1}{m}}^1 \frac{dt_0}{t_0} \prod_{\mu=1}^{n-1} \sum_{w_\mu=0}^{\infty} \frac{(1-t_0)^{s+n-1+\sum_{d=1}^{n-1} w_d}}{s+\mu+\sum_{d=1}^{\mu} w_d} \\
&= n! \prod_{\mu=0}^n \sum_{w_\mu=\delta_{\mu 0}}^{\infty} \frac{\left(1-\frac{1}{m}\right)^{n+\sum_{d=0}^n w_d}}{\mu+\sum_{d=0}^{\mu} w_d}
\end{aligned} \tag{6}$$

where $\delta_{\mu 0}$ is the Kronecker Delta function. This is “simplified” by considering the Sterling Numbers of the First Kind explicitly - however, I will stick with this definition for now. Alternatively, to series expand the other term in Equation 5 is far simpler:

$$t^m = \sum_{\ell=0}^{\infty} \frac{m^{\ell} \ln^{\ell}(t)}{\ell!} \quad (7)$$

Now, revisiting Equation 5, we have:

$$\begin{aligned} \gamma_n = \lim_{m \rightarrow \infty} -n! \prod_{\mu=0}^n \sum_{w_{\mu}=\delta_{\mu 0}}^{\infty} \frac{(1 - \frac{1}{m})^{n + \sum_{d=0}^n w_d}}{\mu + \sum_{d=0}^{\mu} w_d} - \int_0^1 dt \\ \sum_{\omega=0}^n \left(\binom{n}{\omega} \sum_{\ell=1}^{\infty} \frac{m^{\ell} \ln^{\ell}(t)}{(1-t)\ell!} (-\ln(-\ln(t)))^{n-\omega} B_{\omega} \left((-1)^0 \psi_0(1), \dots, (-1)^{n-1} \psi_{n-1}(1) \right) \right) \end{aligned} \quad (8)$$

Now attempting the integral, let $u = -\ln(t)$, thus $dt = e^{-u} du$. Thus:

$$\begin{aligned} \int_0^1 dt \frac{\ln^{\ell}(t) \ln^{n-\omega}(-\ln(t))}{1-t} &= (-1)^{\ell} \int_0^{\infty} du \frac{e^{-u} u^{\ell} \ln^{n-\omega}(u)}{1-e^{-u}} \\ &= (-1)^{\ell} \sum_{\sigma=0}^{\infty} \int_0^{\infty} du e^{-(\sigma+1)u} u^{\ell} \ln^{n-\omega}(u) \\ &= (-1)^{\ell} \sum_{\sigma=0}^{\infty} \int_{-\infty}^{\infty} d\chi e^{-(\sigma+1)e^{\chi}} e^{(\ell+1)\chi} \chi^{n-\omega} \end{aligned} \quad (9)$$

The following remain open questions:

- how to evaluate the integral in Equation 9?
- how to evaluate the sums in ω of the complete exponential Bell Polynomials?
- how to cancel the divergent parts of the two terms of Equation 8 in terms of powers of m ?

Indeed, understanding the first two question will greatly help with understanding the third question. In particular, the integral in terms of t will be the topic of immediate future investigation. Another potential point of improvement is to evaluate if $\ln^{n+1}(m)$ can be written as a first order simplification in terms of $\frac{1}{m}$. All to come in the future. Happy Canada Day 2026! This concludes this blog post.