

# **Title:** Revisiting the Complex Fourier Series

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Recall the formulation of the complex Fourier Series:

$$\begin{aligned} g(x) &= \sum_{n=-\infty}^{\infty} c_n e^{2\pi i n \frac{x}{L}}, \\ c_n &= \frac{1}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} dx' g(x') e^{-2\pi i n \frac{x'}{L}} \end{aligned} \quad (1)$$

This formulation is independent of  $L$  given that  $\frac{L}{2} > x$ . This means that we ought to be able to integrate the top line of Equation 1 in  $L$  without changing the answer. This statement is expressed in the following equation:

$$g(x) = \frac{1}{L - 2|x|} \int_{2|x|}^L dL' \sum_{n=-\infty}^{\infty} \int_{-\frac{L'}{2}}^{\frac{L'}{2}} dx' \frac{g(x') e^{2\pi i n \frac{x-x'}{L'}}}{L'} \quad (2)$$

Switching the order of integration in  $L'$  and  $x'$ , we find:

$$g(x) = \sum_{n=-\infty}^{\infty} \left( \int_{-|x|}^{|x|} dx' \int_{2|x|}^L dL' + \int_{|x|}^{\frac{L}{2}} dx' \int_{2x'}^L dL' + \int_{-\frac{L}{2}}^{-|x|} dx' \int_{-2x'}^L dL' \right) \frac{g(x') e^{2\pi i n \frac{x-x'}{L'}}}{(L - 2|x|) L'} \quad (3)$$

Now we are left with integrals of the form:

$$I = \int_a^b dt \frac{e^{\frac{1}{t}}}{t} \quad (4)$$

Substituting  $u = -\frac{1}{t}$ , we get:

$$I = - \int_{-\frac{1}{a}}^{-\frac{1}{b}} du \frac{e^{-u}}{u} = -\text{Ei}\left(\frac{1}{b}\right) + \text{Ei}\left(\frac{1}{a}\right) \quad (5)$$

where Ei is the exponential integral, defined as:

$$\text{Ei}(x) = - \int_{-x}^{\infty} dt \frac{e^{-t}}{t} \quad (6)$$

where we note that this integral must be understood in terms of the Cauchy principal value since the integrand diverges at  $t \rightarrow 0$ . In particular, the following fact allows the integral to be defined:

$$\begin{aligned} \lim_{t \rightarrow 0^-} \frac{e^{-t}}{t} &\rightarrow -\infty \\ \lim_{t \rightarrow 0^+} \frac{e^{-t}}{t} &\rightarrow +\infty \end{aligned} \quad (7)$$

Thus, the integral diverges oppositely on each side of  $t \rightarrow 0$ , such that the diverging parts cancel and only the finite parts contribute to the value of  $\text{Ei}(x)$ . Proceeding onwards, using Equation 5 to compute Equation 3, we get:

$$\begin{aligned}
g(x) = & \sum_{n=-\infty}^{\infty} \int_{-|x|}^{|x|} dx' \frac{g(x')}{(L - 2|x|)} \left( \text{Ei}\left(\frac{2\pi i n(x-x')}{2|x|}\right) - \text{Ei}\left(\frac{2\pi i n(x-x')}{L}\right) \right) \\
& + \int_{|x|}^{\frac{L}{2}} dx' \frac{g(x')}{(L - 2|x|)} \left( \text{Ei}\left(\frac{2\pi i n(x-x')}{2x'}\right) - \text{Ei}\left(\frac{2\pi i n(x-x')}{L}\right) \right) \\
& + \int_{-\frac{L}{2}}^{-|x|} dx' \frac{g(x')}{(L - 2|x|)} \left( \text{Ei}\left(-\frac{2\pi i n(x-x')}{2x'}\right) - \text{Ei}\left(\frac{2\pi i n(x-x')}{L}\right) \right)
\end{aligned} \tag{8}$$

This is an interesting conclusion, as it seems that the Fourier series has been transformed without changing the value of  $g(x)$ , and  $L$  is still a free parameter in the equation. It would seem natural to analyze the equation by taking the limit as  $L \rightarrow \infty$ , but this proves difficult, as one must account for how the ratio of  $\frac{n}{L}$  evolves as one sums in  $n$  towards  $\pm\infty$ , as this is a ratio of two different variables both limiting towards  $\infty$ . As well, the integrals in  $x'$  have upper and lower bounds that involve  $L$ , meaning that the bounds also go to  $\pm\infty$ .

There are other questions that are pertinent. From previous investigations on this blog, the Fourier series prior to integrating in  $L$  has the following form:

$$g(x) = \lim_{N \rightarrow \infty} \int_{-\frac{L}{2}}^{\frac{L}{2}} dx' g(x') \left[ \frac{\sin\left(2\pi N \frac{x-x'}{L}\right) \cot\left(\pi \frac{x-x'}{L}\right)}{L} \right] \tag{9}$$

The previous discussion outlined the idea that this function behaves like a Dirac Delta function, but is not a Dirac delta function in that it is not 0 everywhere and infinite at some value, but instead is an infinite frequency sine wave with a divergence at  $x' = x$  such that when integrating, all values average to 0 (by the sine wave) but the  $x' = x$  value which is infinite and does not average out. Averaging this function in  $L$  as we have done above should not change this behaviour, particularly as long as  $N \gg L$ . It is often stated that taking the  $L$  in a Fourier Series to infinity leads to the Fourier transform - this would be an interesting idea to explore in this alternative formulation, though it will not be discussed here.

It would be worthwhile to evaluate Equation 8 for a test function, such as  $g(x) = x$ . We evaluate the integrals and set  $x \rightarrow \frac{\pi}{2}$  and  $L \rightarrow 2\pi$ . We then get an intriguing expression for  $\pi$ :

$$\begin{aligned}
\pi = & \sum_{n \neq 0}^{(-\infty, \infty)} \frac{(-1)^n}{2n} \left( i^n (4i - 2n\pi) + (-1)^n (n\pi - i) + i n^2 \pi^2 \left( \text{Ei}\left(\frac{i n \pi}{2}\right) - \text{Ei}(i n \pi) \right) \right) \\
1 = & \sum_{n \neq 0}^{(-\infty, \infty)} \frac{(-1)^n}{2n\pi} \left( i^n (4i - 2n\pi) + (-1)^n (n\pi - i) + i n^2 \pi^2 \left( \text{Ei}\left(\frac{i n \pi}{2}\right) - \text{Ei}(i n \pi) \right) \right)
\end{aligned} \tag{10}$$

Many other expressions are possible for  $g(x) = x$ , and countless more are possible for different  $g(x)$  functions. It may also be possible to integrate Equation 8 again in terms of  $L$  and apply the same formalism to obtain more alternative Fourier Series expressions. This concludes this blog post.